

## A Simplified Kinetic Model for the Tropospheric Ozone Cycle

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**ABSTRACT.** To better understand the dynamics of air pollutants, several mathematical or computational models have been developed and employed. Among the pollutants of interest are those related to the cycle of formation of tropospheric ozone ( $O_3$ ), which involve nitrogen oxides  $NO_x$ . In this process, there is a sequence of chemical reactions whose most elementary modeling can be described in terms of ordinary differential equations (ODEs), in which the concentrations (in  $\mu\text{g}/\text{m}^3$ ) of gases ( $O$ ,  $NO$ ,  $NO_2$ , and  $O_3$ ) are functions of time. A novel study of the model is presented in terms of the qualitative theory of ordinary differential equations, for which the steady state of interest is non-hyperbolic. To study it, the Center Manifold Theorem was used to determine its stability. As for the results, our analytical calculations demonstrate the asymptotic local stability of the steady state, which was also numerically corroborated. Other than this new result of stability, the conclusion is that the simplified model of ozone kinetics with fixed kinetic parameters does not allow the behavior of sustained oscillatory solutions for the referred concentrations of pollutants, requiring other ingredients for this to be feasible.

**Keywords:** mathematical modeling, ordinary differential equations, air pollutants, air quality, nitrogen oxides.

### 1 INTRODUCTION

In environmental research and management, there is a clear need to describe how certain environmental systems evolve from certain known conditions. In this sense, the environmental sciences

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agree with ideas from the physical and chemical sciences as to developing and improving quantitative tools that can describe and modulate their phenomena and processes. However, despite sharing this quantitative view, concerning the complexity of their objects of study, environmental sciences are configured differently [28].

In addition to extensive dimensions, also inherent in environmental systems are multiple events of different natures: physical phenomena, interaction of chemical species, and biological processes. Therefore, environmental sciences constitute an interdisciplinary field of knowledge that is developed in the areas of science, technology, engineering, and mathematics [8], and whose central objective concerns the study and management of processes that impact or may impact human health and the environment. Ultimately, this translates into actions and interventions that simultaneously aim to care for human health and preserve the natural environment. Thus, there is a crucial need to consider the future and understand the impact of environmental events that have not (yet) happened, whether natural (e.g., floods and windstorms) or anthropogenic (e.g., deforestation and air pollution).

In the management of environmental processes and problems, it is essential to prepare forecasts of certain quantities of interest. To this end, in the context of air pollution studies, experimental approaches are crucial, especially for dynamically monitoring air quality indexes. However, in terms of predictive studies, analyzing environmental monitoring data is only the first step in preparing air pollution forecasts. Their development requires resorting to theoretical approaches, which are based on mathematical and computational models.<sup>1</sup> These theoretical-computational approaches make it possible to simulate different atmospheric pollution scenarios, without limitations as to scale, time, and financial resources that are required in experimental approaches. From a practical point of view, these characteristics make mathematical and computational models essential for predicting the behaviors of real environmental systems [7].

In terms of air quality, there are several types of air pollutants. These include atmospheric aerosols and gases (e.g.,  $O_3$ ,  $NO_2$ , and  $NO$ ) [21]. According to the emission source, pollutants can also be classified as primary, when directly emitted by sources, or secondary, when generated by physical and chemical processes in the atmosphere itself [21]. In the troposphere, with pollutants emitted, and with the natural presence of oxygen, several reactions occur (in steps) for the final formation of carbon dioxide ( $CO_2$ ). In these steps, several intermediate compounds are formed. The  $NO_x$  acts as catalysts and sunlight as energy (photochemical reaction) of these reactions, and the main product formed is the ozone gas [21].

As for relevance,  $NO_2$  and  $NO$  can be considered the most important molecules in the chemistry of the atmosphere. In the stratosphere (10 to 25 km altitude), the presence of  $O_3$  is beneficial, mainly due to the formation of the ozone layer. However, in the troposphere (0 to 10 km altitude), its presence is toxic to humans [27]. If at concentrations of inadequate levels,  $NO_2$  and  $O_3$  can cause severe respiratory problems [11, 26, 29]. In this regard, the World Health Organization

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<sup>1</sup>According to the definition of the US National Research Council, and not only in the mathematical scope, a model is “a simplification of reality that is constructed to gain insights into select attributes of a particular physical, biological, economic, or social system” [15].

(WHO) states in its most recent report that there is significant evidence that proves the effects of these pollutants on health [30]. In particular,  $O_3$  and one of its precursors ( $NO_2$ ) follow the guidelines of the Air Quality Guidelines established by the World Health Organization, which must be observed to pose less health risk [29]. Recently, such guidelines have been updated after an extensive review of scientific evidence on the effects of exposure to air pollutants. Their main objective is to guide international legislation, focusing on reducing emissions according to the socioeconomic conditions of each country [3, 13]. Table 1 presents a comparison of these guidelines after the 2021 update for air pollutants directly related to the ozone cycle.

Table 1: WHO 2021/2005 air pollution guidelines comparison.

Pollutant	Average time	WHO 2005	WHO 2021	Change
$O_3$ ( $\mu\text{g}/\text{m}^3$ )	Peak Season	N/A	60	Newly introduced
	8-hour	100	100	Unchanged
	Annual	40	10	Decreasing of 75%
$NO_2$ ( $\mu\text{g}/\text{m}^3$ )	24-hour	N/A	25	Newly introduced
	1-hour	200	200	Unchanged

Given the relevance of the topic to human health, several mathematical and computational models have been employed in the study and prediction of concentrations of pollutants in the tropospheric ozone cycle. This area usually favors approaches involving the diffusion and convection of pollutants, in addition to kinetics containing a high number of chemical reactions, variables, and parameters [12]. However, despite being highly recommended for large-scale applications, these complex approaches tend to disregard some basic facts related to fundamental mechanisms in tropospheric pollution [10, 12]. In this work, we address one of these aspects from a mathematical point of view, to determine whether a simplified kinetic model of the tropospheric ozone with fixed kinetic parameters allows the existence of sustained oscillatory solutions for the referred concentrations of the pollutants of the tropospheric ozone cycle. A notable example of experimental observations of these oscillations in real data is presented in Figure 2 of the reference [20].

Given the presented objective, this work is structured as follows. In Section 2, we describe the basic elements that compose a simplified kinetic model of the tropospheric ozone cycle, which therefore is presented without the presence of external sources of pollutants. In Section 3, we analytically demonstrate and numerically confirm that fixed kinetic parameters cannot generate sustained oscillatory variations for the concentrations of the pollutants of the tropospheric ozone cycle described by the mathematical model. Finally, in Section 4 we present the conclusions of the work and its brief prospects.

## 2 MODELING THE TROPOSPHERIC OZONE CYCLE

Nitrogen dioxide ( $NO_2$ ) and nitric oxide (NO) are jointly referred to as ( $NO_x$ ) (nitrogen oxides) and enter the atmosphere through natural processes (biological processes and atmospheric

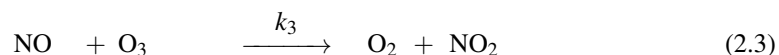
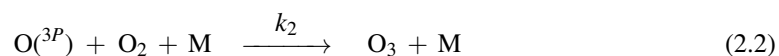
discharges) and anthropogenic processes (combustion of fossil fuels (electricity and transport), biomass burning, and soil emissions from terrestrial ecosystems) [25]. Tropospheric ozone ( $O_3$ ) is formed by a series of photochemical reactions, which involve nitrogen oxides ( $NO_x$ ) and volatile organic compounds (VOCs), which are unburnt or partially oxidized hydrocarbons with boiling temperatures between 50 and 260 °C [2]. Such compounds can be emitted by the burning of fossil fuels (e.g., in vehicle emissions) [14], by the evaporation of fuels, by the use of solvents, by the burning of biomass for energy generation, in forest fires and in fires for agricultural purposes [1, 24].

Chemically, the formation of  $O_3$  occurs through the association of oxygen atoms (O) with the molecular oxygen ( $O_2$ ) present in the atmosphere, having as the main source the dissociation of  $NO_2$  due to sunlight [2]. In turn,  $NO_2$  is formed by the reaction of NO and VOCs with the aid of free radicals [18]. In urban areas, the formation reaction of  $O_3$  is limited by the presence and frequency of NO emission. As these pollutants are transported in air masses to other places, such as rural areas, where there is no significant NO emission, the degradation of the formed and/or transported does not occur, causing measurement peaks in these areas, even though they were not responsible for their generation [20].

From a chemical point of view, the non-linearity in the process of formation of  $O_3$  is due to the sequence of reactions involved. The increases and decreases in ozone concentration occur as a function of the mixing ratio between  $NO_x$  and VOCs. There are two different sensitivity regimes for the formation of this pollutant due to its precursors ( $NO_x$  and VOCs). In the regime sensitive to  $NO_x$ , in which there is a low concentration of  $NO_x$  and a high concentration of VOCs, the formation of  $O_3$  increases with the increase of  $NO_x$ . On the other hand, in the regime sensitive to VOCs, there is a low concentration of VOCs and a high concentration of  $NO_x$ , and the formation of  $O_3$  increases with the decrease of  $NO_x$  [22]. In the atmosphere, in addition to VOCs, there are also other external pollutants formed in incomplete combustion processes that can impact the formation of tropospheric ozone. However, in this study, the tropospheric ozone cycle is analyzed as a closed system, without the contribution of pollutant emission sources.

## 2.1 Kinetics of the Tropospheric Ozone Cycle with no Emission of Pollutants

For the constitution of the compartments of the mathematical model, we considered the concentrations (in  $\mu\text{g}/\text{m}^3$ ) of the following gases over time:  $O(^3P)$ , NO,  $NO_2$ , and  $O_3$ , where  $O(^3P)$  denotes the oxygen atom in its fundamental state. In the troposphere, these gases react according to the following physicochemical reactions [9, 12, 17, 19, 27]:



where reaction (2.1) refers to the photolysis of nitrogen dioxide, reaction (2.2) refers to the formation of ozone, in which the oxygen atom reacts with the molecular oxygen generating ozone, and reaction (2.3) describes the consumption of ozone by reaction with nitric oxide, again forming nitrogen dioxide. In (2.1), (2.2), and (2.3),  $k_i$  ( $i = 1, 2, 3$ ) denotes the kinetic parameters of the referred physicochemical reactions. Since  $k_1, k_2$  and  $k_3$  refer to kinetic rates, these quantities are positive parameters of the mathematical model.

Theoretically, the temporal dynamics of the chemical kinetics given by reactions (2.1), (2.2) and (2.3) can be described by the following system of ordinary differential equations (cf. [9, 17, 19]):

$$\left\{ \begin{array}{l} \frac{d[O(^3P)]}{dt} = k_1[NO_2] - k_2[O(^3P)] \quad (2.4) \\ \frac{d[NO]}{dt} = k_1[NO_2] - k_3[NO][O_3] \quad (2.5) \\ \frac{d[NO_2]}{dt} = k_3[NO][O_3] - k_1[NO_2] \quad (2.6) \\ \frac{d[O_3]}{dt} = k_2[O(^3P)] - k_3[NO][O_3] \quad (2.7) \end{array} \right.$$

where  $k_1$  and  $k_2$  are first-order kinetic parameters and  $k_3$  is a second-order kinetic parameter. Initial condition is non-negative.

Equations (2.4), (2.5), (2.6), and (2.7) describe the variation of the concentration of each of the polluting gases over time. They mathematically model the rate at which reactants become products, where each portion of the direct side of the equations refers to the formation (+) or consumption (-) of the compounds involved. Moreover, in general, since the ozone cycle is photon-guided,  $k_i$  parameters vary in relation to time of day and period of year [27].

In equation (2.4), referring to the concentration of  $O(^3P)$ , there is formation and consumption of  $O(^3P)$ . The formation of  $O(^3P)$  occurs as a result of the  $NO_2$  photodissociation reaction, and its consumption occurs by the  $O_3$  synthesis reaction. In equation (2.5),  $NO$  is formed by  $NO_2$  photodissociation reaction, and consumed by the oxidation of  $NO$  by  $O_3$ . In equation (2.6), in turn, the amount of  $NO_2$  is determined by the rate of oxidation of  $NO$  by  $O_3$  and the rate of photodissociation of  $NO_2$ . Finally, in equation (2.7), the variation in the concentration of  $O_3$  occurs by the consumption of the pollutant, in oxidizing  $NO$ , and by its formation proportional to the concentration of  $O(^3P)$  [17, 27].

Despite being essential for the removal of  $O(^3P)$  and formation of  $O_3$ , the molecular oxygen  $O_2$  involved in the ozone synthesis reaction is not present in the system (2.4)-(2.7). This is mainly because  $O_2$  is abundant in the atmosphere and its concentration does not vary significantly as a function of solar radiation. Therefore, the concentration of  $O_2$  is considered constant over time [17].

### 3 RESULTS

#### 3.1 Existence and Uniqueness of the Solution

The existence and uniqueness of the solution of the system (2.4)-(2.7) is a direct consequence of the Fundamental Existence-Uniqueness Theorem of Picard [16].

#### 3.2 Positivity

To prove that the solution remains positive, for a dynamical system of the form  $X' = f(X)$ , where  $X \in \mathbb{R}^n$  and  $f : \mathbb{R}^n \rightarrow \mathbb{R}^n$  is Lipschitz continuous, one can check that each component  $f_i \geq 0$  if  $X_i = 0$  for all  $i = 1, 2, \dots, n$ , which is known as tangential condition [6]. In fact, this condition is fulfilled for the system (2.4)-(2.7): for each equation of the referred system, the respective right-hand side is always positive when the respective variable is set to zero. For example, in the case of the equation for  $d[\text{O}^{(3P)}]/dt$ , setting  $[\text{O}^{(3P)}]$  equals to zero in the equation (2.4) results in a right-hand side  $k_1 [\text{NO}_2] \geq 0$ .

#### 3.3 Stability Analysis

When analyzing the dynamics of the chemical kinetics of the reactions that characterize the ozone cycle, its possible steady states must also be considered. A steady state is one in which time-dependent variables have constant values that make zero all rates of change (or, in this case, that make zero the reaction rates).

To study the system (2.4)-(2.7) in terms of the qualitative theory of ordinary differential equations, we denote the time-varying concentrations of the chemical species at the moment  $t$  as  $C_1(t) \doteq [\text{O}^{(3P)}](t)$ ,  $C_2(t) \doteq [\text{NO}](t)$ ,  $C_3(t) \doteq [\text{NO}_2](t)$  and  $C_4(t) \doteq [\text{O}_3](t)$ . The equilibrium solutions of (2.4)-(2.7) are those in which the rates of variation of  $[\text{O}^{(3P)}]$ ,  $[\text{NO}]$ ,  $[\text{NO}_2]$ , and  $[\text{O}_3]$  are all null. Therefore, we have the following system of algebraic equations:

$$\begin{cases} k_1 C_3^* - k_2 C_1^* = 0, & (3.1) \\ k_1 C_3^* - k_3 C_2^* C_4^* = 0, & (3.2) \\ k_3 C_2^* C_4^* - k_1 C_3^* = 0, & (3.3) \\ k_2 C_1^* - k_3 C_2^* C_4^* = 0. & (3.4) \end{cases}$$

Not only in (3.1)-(3.4), but also in the system (2.4)-(2.7), there are dependency relations between these equations:

$$\frac{dC_1}{dt} = -\frac{dC_3}{dt} - \frac{dC_4}{dt}, \quad (3.5)$$

which couple the temporal variation rates of the variables  $[\text{O}^{(3P)}]$ ,  $[\text{NO}_2]$  and  $[\text{O}_3]$ , in which the variation of  $\text{NO}_2$  concentration depends on the consumption of  $\text{O}^{(3P)}$  and  $\text{O}_3$ . The other dependency relation is

$$\frac{dC_2}{dt} = -\frac{dC_3}{dt}, \quad (3.6)$$

which indicates that NO formation rate is equal to NO<sub>2</sub> consumption rate. In practice, it is through this redundancy that the concentration of one of these gases is experimentally determined based on the concentration of the other. Thus, by means of equations (3.1) and (3.2), and given that the parameters  $k_i$  ( $i = 1, 2, 3$ ) are non-zero, we have the equalities:

$$C_1^* = \frac{k_1}{k_2} C_3^*, \tag{3.7}$$

and

$$C_2^* C_4^* = \frac{k_1}{k_3} C_3^* \implies C_3^* = \frac{k_3}{k_1} C_2^* C_4^*. \tag{3.8}$$

Replacing  $C_3^*$  in equation (3.7),

$$C_1^* = \frac{k_1 k_3}{k_2 k_1} C_2^* C_4^* = \frac{k_3}{k_2} C_2^* C_4^*. \tag{3.9}$$

Thus, explicitly, the equilibrium solution of the system is given by

$$E = \left( \frac{k_3}{k_2} C_2^* C_4^*, C_2^*, \frac{k_3}{k_1} C_2^* C_4^*, C_4^* \right), \tag{3.10}$$

where the values of  $[NO]^*$  and  $[O_3]^*$  determine the values of  $[O^{(3P)}]^*$  and of  $[NO_2]^*$ . The Jacobian matrix at equilibrium state  $E$  in (3.10) is given by:

$$J(C_1^*, C_2^*, C_3^*, C_4^*) = \begin{pmatrix} -k_2 & 0 & k_1 & 0 \\ 0 & -k_3 C_4^* & k_1 & -k_3 C_2^* \\ 0 & k_3 C_4^* & -k_1 & k_3 C_2^* \\ k_2 & -k_3 C_4^* & 0 & -k_3 C_2^* \end{pmatrix}, \tag{3.11}$$

whose eigenvalues are determined by

$$\det(J - \lambda I) = 0, \tag{3.12}$$

where, by development into cofactors, we obtain:

$$\lambda^2 [k_2 k_3 C_2^* + k_2 k_3 C_4^* + k_1 k_2 + (k_1 + k_2 + k_3 (C_2^* + C_4^*)) \lambda + k_1 k_3 C_2^* + \lambda^2] = 0. \tag{3.13}$$

Thus, the characteristic polynomial (3.13) has four roots, two of which are null:  $\lambda_{1,2} = 0$ . As for non-null eigenvalues, they are roots of the following polynomial equation:

$$a \lambda^2 + b \lambda + c = 0, \tag{3.14}$$

whose coefficients are all positive:  $a \doteq 1$ ,  $b \doteq [k_1 + k_2 + k_3 (C_2^* + C_4^*)]$  and  $c \doteq k_1 k_2 + k_2 k_3 (C_2^* + C_4^*) + k_1 k_3 C_2^*$ . Therefore, we use the solver of the second degree equation to characterize the solutions of the polynomial,  $\Delta = b^2 - 4c$ , given by

$$\Delta = k_1^2 + 2k_1 k_3 C_4^* + k_3^2 (C_4^*)^2 + k_2^2 - 2k_2 k_3 C_2^* + k_3^2 (C_2^*)^2 - 2k_1 k_2 - 2k_1 k_3 C_2^* - 2k_2 k_3 C_4^* + 2k_3^2 C_2^* C_4^* \tag{3.15}$$

Adding and subtracting the element  $4k_1k_3C_2^*$ , we have:

$$\Delta = [(k_1 + k_3C_4^*) - (k_2 - k_3C_2^*)]^2 - 4k_1k_3C_2^*. \tag{3.16}$$

Thus, in addition to the two null eigenvalues, the other two remaining eigenvalues are:

$$\lambda_{3,4} = \frac{-b \pm \sqrt{\Delta}}{2}. \tag{3.17}$$

If  $\Delta = 0$ ,  $\lambda_{3,4} = -b/(2)$ , which occurs if

$$\frac{1}{k_1k_3C_2^*} \left( \frac{(k_1 + k_3C_4^*) - (k_2 - k_3C_2^*)}{2} \right)^2 = 1. \tag{3.18}$$

If  $\Delta < 0$ , the roots are imaginary and complex conjugates, of the form

$$\lambda_{3,4} = \frac{-b \pm i\sqrt{|\Delta|}}{2}, \tag{3.19}$$

which occurs if

$$\frac{1}{k_1k_3C_2^*} \left( \frac{(k_1 + k_3C_4^*) - (k_2 - k_3C_2^*)}{2} \right)^2 < 1. \tag{3.20}$$

In both cases, for  $\Delta = 0$  and for  $\Delta < 0$ , the real part of the non-null roots is always negative by the definition of the parameters that define the coefficient  $b$ .

When  $\Delta > 0$ , as the coefficients  $a, b$  and  $c$  are positive, then  $b^2 - 4c < b^2$ , since  $-4c < 0$  and  $b^2 > 0$ . In this case,

$$\sqrt{b^2 - 4c} < \sqrt{b^2}, \tag{3.21}$$

$$\sqrt{b^2 - 4c} < |b|, \tag{3.22}$$

$$\sqrt{b^2 - 4c} < b, \tag{3.23}$$

$$-b + \sqrt{b^2 - 4c} < 0, \tag{3.24}$$

and, as  $a > 0$ , the other eigenvalues are:

$$\lambda_{3,4} = \frac{-b \pm \sqrt{b^2 - 4c}}{2} < 0, \tag{3.25}$$

which occurs if

$$\frac{1}{k_1k_3C_2^*} \left( \frac{(k_1 + k_3C_4^*) - (k_2 - k_3C_2^*)}{2} \right)^2 > 1. \tag{3.26}$$

To conclude, therefore, we observe that the eigenvalues always have a negative real part or are null, as compiled in Table 2 where we already use the value  $a \doteq 1$  as defined in equation (3.14).



Table 2: Eigenvalues associated with the equilibrium state  $E$ , Equation (3.10).

Discriminant	Condition	Eigenvalues
$\Delta > 0$	$\frac{1}{k_1 k_3 C_2^*} \left( \frac{(k_1 + k_3 C_4^*) - (k_2 - k_3 C_2^*)}{2} \right)^2 > 1$	$\lambda_{1,2} = 0, \lambda_{3,4} = \frac{-b \pm \sqrt{\Delta}}{2}$
$\Delta = 0$	$\frac{1}{k_1 k_3 C_2^*} \left( \frac{(k_1 + k_3 C_4^*) - (k_2 - k_3 C_2^*)}{2} \right)^2 = 1$	$\lambda_{1,2} = 0, \lambda_{3,4} = -\frac{b}{2}$
$\Delta < 0$	$\frac{1}{k_1 k_3 C_2^*} \left( \frac{(k_1 + k_3 C_4^*) - (k_2 - k_3 C_2^*)}{2} \right)^2 < 1$	$\lambda_{1,2} = 0, \lambda_{3,4} = \frac{-b \pm i\sqrt{ \Delta }}{2}$

As shown in Table 2, given the existence of null eigenvalues, the equilibrium state  $E$  of equation (3.10) is non-hyperbolic. To characterize its stability, we use the Center Manifold Theorem [4, 16], denoting the Jacobian of the function  $f$  by  $f'$ .

**Theorem 1.** (Local Center Manifold Theorem). *Let  $f \in C^r(K)$ , where  $r \geq 1$  and  $K$  is an open subset of  $\mathbb{R}^n$  containing the origin. Suppose  $f(0) = 0$  and that  $f'(0)$  has  $c$  eigenvalues with null real part and  $s$  eigenvalues with negative real part, where  $c + s = n$ . The system (2.4)-(2.7) can be written in diagonal form as,*

$$\begin{cases} \frac{dx}{dt} = Cx + F(x, y), & (3.27) \\ \frac{dy}{dt} = Py + G(x, y), & (3.28) \end{cases}$$

where  $(x, y) \in \mathbb{R}^c \times \mathbb{R}^s$ ,  $C$  is a square matrix with  $c$  eigenvalues with zero real part,  $P$  is a square matrix with  $s$  eigenvalues with negative real part,  $F(0) = G(0) = 0$ , and  $F'(0) = G'(0) = 0$ . Thus, there is a  $\delta > 0$  and a function  $h \in C^r(N_\delta(0))$  that defines a local center manifold given by

$$W_{local}^c = \{(x, y) \in \mathbb{R}^c \times \mathbb{R}^s \mid y = h(x) \text{ for } |x| < \delta\}, \tag{3.29}$$

where  $W^c$  is tangent to the central subspace  $E^c = \{(x, y) \in \mathbb{R}^c \times \mathbb{R}^s \mid y = 0\}$  and satisfies

$$h'(x)[Cx + F(x, h(x))] - Ph(x) - G(x, h(x)) = 0, \tag{3.30}$$

for  $|x| < \delta$ , and the flow over the center manifold  $W^c(0)$  is defined by the system of equations

$$\frac{dx}{dt} = Cx + F(x, h(x)), \tag{3.31}$$

$\forall x \in \mathbb{R}^c$ , with  $|x| < \delta$ . To use the Center Manifold Theorem at equilibrium state  $E$  it is necessary to make a translation to the origin. To do so, we use the following change of variables:

$$C_{11} = C_1 - \frac{k_3}{k_2} C_2^* C_4^*, \tag{3.32}$$

$$C_{12} = C_2 - C_2^*, \tag{3.33}$$

$$C_{13} = C_3 - \frac{k_3}{k_1} C_2^* C_4^*, \tag{3.34}$$

$$C_{14} = C_4 - C_4^*. \tag{3.35}$$

Therefore, in these translated variables, the system is given by:

$$dC_{11}/dt = k_1 C_{13} - k_2 C_{11}, \tag{3.36}$$

$$dC_{12}/dt = k_1 C_{13} - k_3 C_{12} C_{14} - k_3 C_4^* C_{12} - k_3 C_4^* C_{14}, \tag{3.37}$$

$$dC_{13}/dt = k_3 C_{12} C_{14} + k_3 C_4^* C_{12} + k_3 C_4^* C_{14} - k_1 C_{13}, \tag{3.38}$$

$$dC_{14}/dt = k_2 C_{11} - k_3 C_{12} C_{14} - k_3 C_4^* C_{12} - k_3 C_2^* C_{14}, \tag{3.39}$$

or in the matrix form:

$$\begin{pmatrix} dC_{11}/dt \\ dC_{12}/dt \\ dC_{13}/dt \\ dC_{14}/dt \end{pmatrix} = \begin{pmatrix} -k_2 & 0 & k_1 & 0 \\ 0 & -k_3 C_4^* & k_1 & -k_3 C_2^* \\ 0 & k_3 C_4^* & -k_1 & k_3 C_2^* \\ k_2 & -k_3 C_4^* & 0 & -k_3 C_2^* \end{pmatrix} \begin{pmatrix} C_{11} \\ C_{12} \\ C_{13} \\ C_{14} \end{pmatrix} + \begin{pmatrix} 0 \\ -k_3 C_{12} C_{14} \\ k_3 C_{12} C_{14} \\ -k_3 C_{12} C_{14} \end{pmatrix}. \tag{3.40}$$

To determine whether the matrix of the system (3.40) is diagonalizable, we can study the geometric multiplicity of the null eigenvalue, which has an algebraic multiplicity equal to 2. To obtain it, we solve the equation  $Av = \lambda v$ , where  $A$  is the matrix of the system (3.40),  $\lambda$  is the null eigenvalue and  $v$  is the eigenvector associated with  $\lambda = 0$ :

$$\begin{pmatrix} -k_2 & 0 & k_1 & 0 \\ 0 & -k_3 C_4^* & k_1 & -k_3 C_2^* \\ 0 & k_3 C_4^* & -k_1 & k_3 C_2^* \\ k_2 & -k_3 C_4^* & 0 & -k_3 C_2^* \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ v_3 \\ v_4 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}. \tag{3.41}$$

Solving it, we have:

$$v = \left( \frac{k_3}{k_2} (C_2^* v_4 + C_4^* v_2), v_2, \frac{k_3}{k_1} (C_2^* v_4 + C_4^* v_2), v_4 \right), \tag{3.42}$$

and thus the eigenvalue  $\lambda = 0$  has geometric and algebraic multiplicity equal to 2. Next, we study the three cases for the discriminant that are shown in Table 2.

**Case I:**  $\lambda_{1,2} = 0$  and  $\lambda_{3,4} = (-b \pm \sqrt{\Delta})/2$ , and the matrix of the system (3.40) is diagonalizable. Thus, there is a basis of vectors that forms the matrix  $Q$ , and a diagonal matrix  $D$  with the eigenvalues of  $A$ , such that  $Q^{-1}AQ = D$ , and the system to be studied is:

$$\begin{pmatrix} du_1/dt \\ du_2/dt \\ du_3/dt \\ du_4/dt \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & (-b - \sqrt{\Delta})/2 & 0 \\ 0 & 0 & 0 & (-b + \sqrt{\Delta})/2 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{pmatrix} + \frac{Bk_3}{2\sqrt{\Delta}} \begin{pmatrix} 0 \\ 0 \\ (-b_2 + \sqrt{\Delta}) \\ (+b_2 + \sqrt{\Delta}) \end{pmatrix}. \tag{3.43}$$

where  $b_2 = k_1 + k_2 - (C_2^* + C_4^*)k_3$ ,  $b_1 = -k_1 + k_2 + k_3(C_2^* - C_4^*)$ , and

$$B = \left[ \frac{k_1 u_2}{C_4^* k_3} - \frac{C_2^* u_1}{C_4^*} + 2C_2^* k_3 \left( \frac{u_3}{b_1 + \sqrt{\Delta}} - \frac{u_4}{-b_1 + \sqrt{\Delta}} \right) \right] (u_1 + u_3 + u_4)$$

and  $u_1, u_2, u_3, u_4$  are such that

$$\begin{pmatrix} C_{11} \\ C_{12} \\ C_{13} \\ C_{14} \end{pmatrix} = Q \begin{pmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{pmatrix}.$$

The system (3.43) has a suitable structure to apply Theorem 1. Let  $C = 0$  and

$$P = \begin{pmatrix} \frac{-b - \sqrt{\Delta}}{2} & 0 \\ 0 & \frac{-b + \sqrt{\Delta}}{2} \end{pmatrix}, \tag{3.44}$$

where we have already shown that  $(-b - \sqrt{\Delta})/2$  and  $(-b + \sqrt{\Delta})/2$  are negative,  $F = (F_1, F_2)$ , where  $F_1(u_1, u_2, u_3, u_4) = 0$  and  $F_2(u_1, u_2, u_3, u_4) = 0$ ,  $G = (G_1, G_2)$ , where  $G_1(u_1, u_2, u_3, u_4) = Bk_3(-b_2 + \sqrt{\Delta})/(2\sqrt{\Delta})$ , and  $G_2(u_1, u_2, u_3, u_4) = Bk_3(+b_2 + \sqrt{\Delta})/(2\sqrt{\Delta})$ . Therefore,  $F(0, 0, 0, 0) = 0$ ,  $F'(0, 0, 0, 0) = 0$ ,  $G(0, 0, 0, 0) = 0$  and  $G'(0, 0, 0, 0) = 0$ .

Defining  $u_3 = \phi_1(u_1, u_2)$  and  $u_4 = \phi_2(u_1, u_2)$ , the center manifold of dimension two has the following form:

$$M_1[\phi_1, \phi_2] = \frac{-(-b - \sqrt{\Delta})\phi_1}{2} - k_3 \frac{\left( \frac{k_1 u_2}{C_4^* k_3} - \frac{C_2^* u_1}{C_4^*} + \frac{2C_2^* k_3 \phi_1}{b_1 + \sqrt{\Delta}} - \frac{2C_2^* k_3 \phi_2}{-b_1 + \sqrt{\Delta}} \right) (u_1 + \phi_1 + \phi_2) (-b_2 + \sqrt{\Delta})}{2\sqrt{\Delta}} = 0, \tag{3.45}$$

$$M_2[\phi_1, \phi_2] = \frac{-(-b + \sqrt{\Delta})\phi_2}{2} - k_3 \frac{\left( \frac{k_1 u_2}{C_4^* k_3} - \frac{C_2^* u_1}{C_4^*} + \frac{2C_2^* k_3 \phi_1}{b_1 + \sqrt{\Delta}} - \frac{2C_2^* k_3 \phi_2}{-b_1 + \sqrt{\Delta}} \right) (u_1 + \phi_1 + \phi_2) (+b_2 + \sqrt{\Delta})}{2\sqrt{\Delta}} = 0, \tag{3.46}$$

In general, from a practical point of view, the following approximation theorem is sufficient to prove stability [4].

**Theorem 2.** Let  $\phi : \mathbb{R}^n \rightarrow \mathbb{R}^n$  of class  $C^1$  satisfying  $\phi(0) = \nabla\phi(0) = 0$ . Let  $M[\phi](X) = O(|x|^q)$ ,  $x \rightarrow 0$ , for some  $q > 1$ . Thus, there is a center manifold  $\sigma$  satisfying  $\sigma(x) - \phi(x) = O(|x|^q)$ ,  $x \rightarrow 0$ .

By Theorem 1, the equation does not depend on  $h$  and the stability of the system (3.40) is governed by:

$$\begin{pmatrix} du_1/dt \\ du_2/dt \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}. \tag{3.47}$$

Thus, using Theorem 1, the equilibrium state  $E$  is stable for  $\Delta > 0$ .

**Case II:**  $\lambda_{1,2} = 0$ ,  $\lambda_3 = (-b + i\sqrt{|\Delta|})/2$  and  $\lambda_4 = (-b - i\sqrt{|\Delta|})/2$ . In this case, the matrix of the system is also diagonalizable and the system to be studied is:

$$\begin{pmatrix} du_1 \\ du_2 \\ du_3 \\ du_4 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & (-b+i\sqrt{|\Delta|})/2 & 0 \\ 0 & 0 & 0 & (-b-i\sqrt{|\Delta|})/2 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{pmatrix} + \frac{Bk_3}{2i\sqrt{|\Delta|}} \begin{pmatrix} 0 \\ 0 \\ (-b_2+i\sqrt{|\Delta|}) \\ (+b_2+i\sqrt{|\Delta|}) \end{pmatrix}. \tag{3.48}$$

The system (3.48) is written in a suitable structure to apply the Theorem 1.

Let  $C = 0$  and

$$P = \begin{pmatrix} \frac{-b+i\sqrt{|\Delta|}}{2} & 0 \\ 0 & \frac{-b-i\sqrt{|\Delta|}}{2} \end{pmatrix}, \tag{3.49}$$

where we have already shown that  $(-b - i\sqrt{|\Delta|})/2$  and  $(-b + i\sqrt{|\Delta|})/2$  have a negative real part,  $F = (F_1, F_2)$ , where  $F_1(u_1, u_2, u_3, u_4) = 0$  and  $F_2(u_1, u_2, u_3, u_4) = 0$ ,  $G = (G_1, G_2)$ , where  $G_1(u_1, u_2, u_3, u_4) = Bk_3(-b_2 + i\sqrt{|\Delta|})/(2i\sqrt{|\Delta|})$ , and  $G_2(u_1, u_2, u_3, u_4) = Bk_3(+b_2 + i\sqrt{|\Delta|})/(2i\sqrt{|\Delta|})$ . Therefore,  $F(0, 0, 0, 0) = 0$ ,  $F'(0, 0, 0, 0) = 0$ ,  $G(0, 0, 0, 0) = 0$  and  $G'(0, 0, 0, 0) = 0$ .

Thus, defining  $u_3 = \phi_1(u_1, u_2)$  and  $u_4 = \phi_2(u_1, u_2)$ , the center manifold has the following form:

$$M_1[\phi_1, \phi_2] = \frac{-(-b+i\sqrt{|\Delta|})\phi_1}{2} - k_3 \frac{\left(\frac{k_1 u_2}{C_4^* k_3} - \frac{C_2^* u_1}{C_4^*} + \frac{2C_2^* k_3 \phi_1}{b_1 + i\sqrt{|\Delta|}} - \frac{2C_2^* k_3 \phi_2}{-b_1 + i\sqrt{|\Delta|}}\right)(u_1 + \phi_1 + \phi_2)(-b_2 + i\sqrt{|\Delta|})}{2i\sqrt{|\Delta|}} = 0, \tag{3.50}$$

$$M_2[\phi_1, \phi_2] = \frac{-(-b-i\sqrt{|\Delta|})\phi_2}{2} - k_3 \frac{\left(\frac{k_1 u_2}{C_4^* k_3} - \frac{C_2^* u_1}{C_4^*} + \frac{2C_2^* k_3 \phi_1}{b_1 + i\sqrt{|\Delta|}} - \frac{2C_2^* k_3 \phi_2}{-b_1 + i\sqrt{|\Delta|}}\right)(u_1 + \phi_1 + \phi_2)(+b_2 + i\sqrt{|\Delta|})}{2i\sqrt{|\Delta|}} = 0, \tag{3.51}$$

By Theorem 1, the equation governing the solution of the system (3.40) is:

$$\begin{pmatrix} du_1/dt \\ du_2/dt \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}. \tag{3.52}$$

Thus, the stability of the system at the origin. Therefore, by Theorem 1, the equilibrium state  $E$  is also stable for  $\Delta < 0$ .

Summarizing, except in the case where the discriminant  $\Delta$  is zero (see Table 2), we have proved that the equilibrium state  $E$  given in equation (3.10) is locally asymptotically stable.

### 3.4 Numerical Results

To illustrate the analytical results practically, the system of ordinary differential equations of the tropospheric ozone cycle was also solved numerically. To this end, a computational implementation was carried out for the numerical solution of the system of ordinary differential equations(2.4), (2.5), (2.6) and (2.7) The system was treated as an initial value problem in which the initial concentration of  $O(^3P)$ , NO, NO<sub>2</sub>, and O<sub>3</sub> was chosen as being zero, 10<sup>12</sup>, 10<sup>10</sup>, and zero molecules/cm<sup>3</sup>, respectively. These values are indicated in the reference [9], Problem 12.3, Page

443. Moreover, still based on this reference (Table 12.3, Page 443), for a temperature  $T = 298$  K, the values assigned for the kinetic parameters were:  $k_1 = 1.7 \cdot 10^{-2} \text{ s}^{-1}$ ,  $k_2 = 1.4 \cdot 10^3 \cdot e^{\left(\frac{1175}{T}\right)}$   $\text{cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$  and  $k_3 = 1.8 \cdot 10^{-12} \cdot e^{\left(\frac{-1370}{T}\right)}$   $\text{cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$ . From the point of view of the numerical solution, the differences of several orders of magnitude in the values of these parameters present a characteristic known as (*stiffness*), which requires a careful and more sophisticated choice than the 4th-order Runge-Kutta method for the realization of the numerical solution of the model [5, 23].

To this end, we used the ode routine of the R package `deSolve`, which can identify and numerically solve ODE (*stiff*) systems. Based on the graph shown in Figure 1, as  $\text{NO}_2$  is consumed,  $\text{NO}$  and  $\text{O}_3$  are formed. In addition, the equilibrium value is reached quickly, approximately at  $t = 200$  seconds. Furthermore,  $\text{NO}_2$  and  $\text{O}_3$  concentrations approximate when stability is reached, where  $\text{NO}_2$  consumption is equivalent to the  $\text{O}_3$  formed. In other words, to have ozone accumulation in the troposphere, other oxidants must participate in the reactions. The curves shown in Figure 1 are the same as the reference cited above (Figure 12.5, Page 443, [9]).

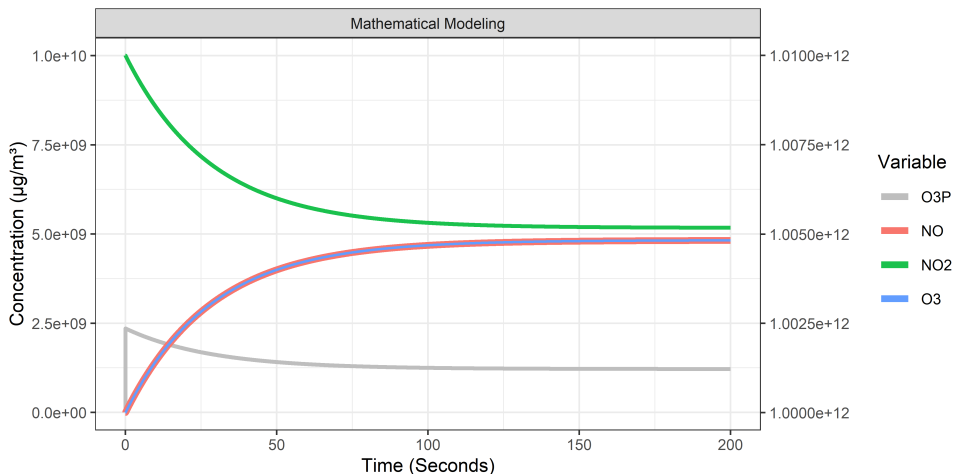


Figure 1: Numerical solution of the mathematical model presented and defined by equations (2.4), (2.5), (2.6), and (2.7) for the tropospheric ozone cycle. The scale of  $[\text{NO}]$  is the vertical axis on the right, and the value of  $[\text{O}^{(3P)}]$  is presented with a multiplicative factor equal to  $10^6$ .

Finally, using the initial condition and the parameter values, we obtained the numerical stationary solution shown in Figure 1:  $(1.2194 \cdot 10^3, 1.0048 \cdot 10^{12}, 5.1786 \cdot 10^9, 4.8213 \cdot 10^9)$ , and the equilibrium solution given by (3.10) is  $(1.2173 \cdot 10^3, 1.0048 \cdot 10^{12}, 5.1699 \cdot 10^9, 4.8213 \cdot 10^9)$ . Using Table 2, and with the referred inequality being less than the unity, the eigenvalues are  $\lambda_{1,2} = 0$ ,  $\lambda_3 = -10573.3$ , and  $\lambda_4 = -61625.7$ . According to the theoretical results, this equilibrium solution is stable. Thus, with this computational simulation, we can graphically verify the fulfillment of the theoretical result that was obtained analytically.

## 4 CONCLUSION

In this work, we present a qualitative study of a mathematical model of ordinary differential equations in time that characterizes the kinetics of the chemical reactions of the ozone cycle in the troposphere. The only equilibrium state of the system is non hyperbolic, and so we used the Center Manifold Theorem to study its local stability. We present complete analytical results that establish that the steady state is stable for any values of the kinetic parameters, which was also corroborated by the numerical solution of the resulting stiff problem. In addition to establishing this unprecedented result of stability, another conclusion is that the simplified model of ozone kinetics with fixed kinetic parameters does not allow the behavior of sustained oscillatory solutions for the referred concentrations of pollutants, requiring other ingredients for this to be feasible. In future works, we intend to address the problem of estimating mathematical model parameters from real data.

### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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